

Compressible MHD Turbulence

Theoretical and astrophysical interests

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Basic Theoretical Assumptions

- Compressible MHD fluid with isothermal closure ($P = C_s^2 \rho$)
- Statistical Homogeneity
- Existence of a stationary state corresponding to the total energy (kinetic + magnetic + compressible) conservation
- Existence of an inertial zone (independent of the external forcing and the dissipation)

Exact relation for isothermal CMHD turbulence (*S.Banerjee & S.Galtier PRE, 2013*)

In the inertial zone & without assuming isotropy

$$\begin{aligned}
 -2\varepsilon &= \frac{1}{2}\nabla_{\mathbf{r}} \cdot \left\langle \left[\frac{1}{2}\delta(\rho\mathbf{z}^-) \cdot \delta\mathbf{z}^- + \delta\rho\delta\mathbf{e} \right] \delta\mathbf{z}^+ + \left[\frac{1}{2}\delta(\rho\mathbf{z}^+) \cdot \delta\mathbf{z}^+ + \delta\rho\delta\mathbf{e} \right] \delta\mathbf{z}^- + \bar{\delta}\left(\mathbf{e} + \frac{v_A^2}{2}\right)\delta(\rho\mathbf{z}^- + \rho\mathbf{z}^+) \right\rangle \\
 &- \frac{1}{8} \left\langle \frac{1}{\beta'}\nabla' \cdot (\rho\mathbf{z}^+ \mathbf{e}') + \frac{1}{\beta}\nabla \cdot (\rho'\mathbf{z}'^+ \mathbf{e}) + \frac{1}{\beta'}\nabla' \cdot (\rho\mathbf{z}^- \mathbf{e}') + \frac{1}{\beta}\nabla \cdot (\rho'\mathbf{z}'^- \mathbf{e}) \right\rangle \\
 &+ \left\langle (\nabla \cdot \mathbf{v}) \left[R'_E - E' - \frac{\bar{\delta}\rho}{2}(\mathbf{v}_{\mathbf{A}'} \cdot \mathbf{v}_{\mathbf{A}}) - \frac{P'}{2} + \frac{P'_M}{2} \right] \right\rangle + \left\langle (\nabla' \cdot \mathbf{v}') \left[R_E - E - \frac{\bar{\delta}\rho}{2}(\mathbf{v}_{\mathbf{A}} \cdot \mathbf{v}_{\mathbf{A}'}) - \frac{P}{2} + \frac{P_M}{2} \right] \right\rangle \\
 &+ \left\langle (\nabla \cdot \mathbf{v}_{\mathbf{A}}) \left[R_H - R'_H + H' - \bar{\delta}\rho(\mathbf{v}' \cdot \mathbf{v}_{\mathbf{A}}) \right] \right\rangle + \left\langle (\nabla' \cdot \mathbf{v}_{\mathbf{A}'}) \left[R'_H - R_H + H - \bar{\delta}\rho(\mathbf{v} \cdot \mathbf{v}_{\mathbf{A}'}) \right] \right\rangle
 \end{aligned}$$

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In the inertial zone & without assuming isotropy

$$\begin{aligned}
 -2\varepsilon &= \frac{1}{2} \nabla_{\mathbf{r}} \cdot \left\langle \left[\frac{1}{2} \delta(\rho \mathbf{z}^-) \cdot \delta \mathbf{z}^- + \delta \rho \delta \mathbf{e} \right] \delta \mathbf{z}^+ + \left[\frac{1}{2} \delta(\rho \mathbf{z}^+) \cdot \delta \mathbf{z}^+ + \delta \rho \delta \mathbf{e} \right] \delta \mathbf{z}^- + \bar{\delta} \left(\mathbf{e} + \frac{v_A^2}{2} \right) \delta(\rho \mathbf{z}^- + \rho \mathbf{z}^+) \right\rangle \\
 &- \frac{1}{8} \left\langle \frac{1}{\beta'} \nabla' \cdot (\rho \mathbf{z}^+ \mathbf{e}') + \frac{1}{\beta} \nabla \cdot (\rho' \mathbf{z}'^+ \mathbf{e}) + \frac{1}{\beta'} \nabla' \cdot (\rho \mathbf{z}^- \mathbf{e}') + \frac{1}{\beta} \nabla \cdot (\rho' \mathbf{z}'^- \mathbf{e}) \right\rangle \\
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 \end{aligned}$$

Reduced form in the presence of a strong B_0

$$-2\varepsilon \simeq \frac{B_0^2}{2} \nabla_{r_{\perp}} \cdot \left\langle \delta \left(\frac{1}{\sqrt{\rho}} \right) \delta(\sqrt{\rho}) \delta \mathbf{v}_{\perp} \right\rangle - \frac{B_0^2}{4} \left\langle (\nabla_{\perp} \cdot \mathbf{v}_{\perp}) \left(1 + \sqrt{\frac{\rho}{\rho'}} \right) + (\nabla'_{\perp} \cdot \mathbf{v}'_{\perp}) \left(1 + \sqrt{\frac{\rho'}{\rho}} \right) \right\rangle$$

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- 2 types of flux terms and 2 types of source terms

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- Unlike incompressible MHD, we do not have separate exact relations for individual pseudo energies

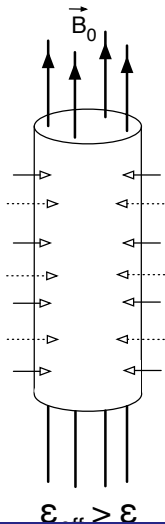
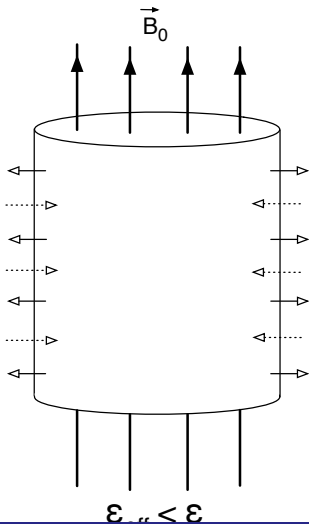
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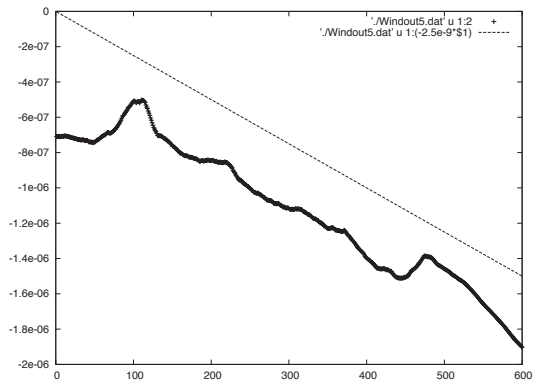
- 2 types of flux terms and 2 types of source terms
- Unlike incompressible MHD, we do not have separate exact relations for individual pseudo energies
- For weak compressibility, 3rd order moment scaling can be expected taking $(\rho z^{+2} z^{-} + \rho z^{-2} z^{+})^{1/3}$ as the scaling variable

Phenomenology for Reduced Case



Scaling in solar wind data

Scaling with 1 day of WIND data with 1 minute plasma resolution



Questions

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- Can we have a phenomenological view for the general case of compressible MHD turbulence ?