

Energy Loss Under Guiding Center Motion

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Breakout Session D:

Pickup Ions and Building the Energetic Particle Reservoir

Basic theory-relativistic energy loss rate [Northrop, 1963]

$$dT/dt = ZeV_{||}E_{||} + (mv_{\perp}^2/2)(\partial \ln B / \partial t + \mathbf{u}_E \cdot \nabla \ln B) + mv_{||}^2 \mathbf{u}_E \cdot (\mathbf{b} \cdot \nabla \mathbf{b})$$

$(mv_{\perp}^2/2)(\partial \ln B / \partial t)$ is betatron term

$\mathbf{u}_E = \mathbf{E} \times \mathbf{B} / B^2$ and $\mathbf{b} \cdot \nabla \mathbf{b} = \boldsymbol{\kappa}$ so remaining terms are

$$\mathbf{u}_E \cdot (mv_{\perp}^2/2 \nabla \ln B + mv_{||}^2 \boldsymbol{\kappa}) = (1/B) \mathbf{E} \cdot [\mathbf{b} \times (mv_{\perp}^2/2 \nabla \ln B + mv_{||}^2 \boldsymbol{\kappa})]$$

Relativistic gradient and curvature drift velocities are

$$\mathbf{u}_G = (M/\gamma ZeB) \mathbf{b} \times \nabla B \quad \mathbf{u}_C = (p_{||}^2/\gamma Z e m_0 B) \mathbf{b} \times \boldsymbol{\kappa}$$

relativistic magnetic moment $M = p_{\perp}^2/2m_0$ p. 29

Terms equal $Ze \mathbf{E} \cdot (\mathbf{u}_G + \mathbf{u}_C)$: *work done on drifting particle by the electric field.*

Momentum fractional time rate of change

$$dT = v dp \text{ so } dT/dt = v dp/dt$$

$$dp/dt = Ze(V_{||}/v)E_{||} + p(1-\mu^2)/2(\partial \ln B / \partial t + \mathbf{u}_E \cdot \nabla \ln B) + p\mu^2 \mathbf{u}_E \cdot (\mathbf{b} \cdot \nabla \mathbf{b})$$

Neglecting the $E_{||}$ term and setting $\mathbf{u}_E = \mathbf{V}_{\perp}$ and $\mathbf{b} \cdot \nabla \mathbf{b} = \boldsymbol{\kappa}$, we have the result [Roelof, 1999; 2000]

$$d \ln p / dt = (1-\mu^2)/2(\partial \ln B / \partial t + \mathbf{V}_{\perp} \cdot \nabla \ln B) + \mu^2 \mathbf{V}_{\perp} \cdot \boldsymbol{\kappa}$$

Remark: The fractional rate is **independent** of particle mass or charge! It depends only on the particle pitch-angle, the component of the transverse plasma velocity and the local spatial-temporal variations in the magnetic field.

Parker field

$$B_r = B_0(r/r_0)^2 \quad B_\Theta = 0 \quad B_\Phi = -B_r(\Omega r \sin\Theta/V) = -B_r \tan\psi$$

General identities

$$\nabla_\perp \ln B = (1/B^2) \mathbf{B} \times (\nabla \times \mathbf{B}) + \boldsymbol{\kappa} \quad \mu_0 \mathbf{J} = \nabla \times \mathbf{B}$$

$$\mu_0 \mathbf{J} = -2(\Omega/V) B_0 (r_0/r)^2 \cos\Theta \mathbf{u}_r$$

$$\mu_0 \mathbf{J} \times \mathbf{B} = -(\Omega B_0/V)^2 (r_0^4/r^3) \sin 2\Theta \mathbf{u}_\Theta$$

thus $\mathbf{V}_\perp \cdot \mathbf{J} \times \mathbf{B} = 0$ and so $\mathbf{V}_\perp \cdot \nabla_\perp \ln B = \mathbf{V}_\perp \cdot \boldsymbol{\kappa}$ for the Parker field

$$\mathbf{V}_{\perp} = (\mathbf{I} - \mathbf{bb}) \cdot \mathbf{V} = V(\mathbf{u}_r \sin^2 \psi + \mathbf{u}_{\phi} \sin \psi \cos \psi)$$

$$B^2 = B_0^2 (r_0/r)^4 (1 + \tan^2 \psi) \quad \partial \tan \psi / \partial \ln r = \tan \psi$$

$$\mathbf{V}_{\perp} \cdot \nabla \ln B = -(V/r)(1 - \cos^4 \psi) = \kappa \cdot \mathbf{V}_{\perp}$$

$$\begin{aligned} d \ln p / dt &= (1 - \mu^2) / 2 (\partial \ln B / \partial t + \mathbf{V}_{\perp} \cdot \nabla \ln B) + \mu^2 \mathbf{V}_{\perp} \cdot \kappa \\ &= (1 - \mu^2) / 2 \partial \ln B / \partial t - (1 + \mu^2) (V / 2r) (1 - \cos^4 \psi) \end{aligned}$$

[Roelof, Proc. 7th Int. Astrophys. Conf., *AIP*, **1039**, 174-183, 2008]

Limit: $\tan \psi \gg 1$ or $r \sin \Theta \gg V / \Omega \sim 1 \text{ AU}$

Also $\partial \ln B / \partial t = 0$, and isotropic p.a.d.

$\langle d \ln p / dt \rangle \rightarrow -2V / 3r$ **“adiabatic cooling”**

General expression for fractional momentum rate

$$d \ln p / dt = (1 - \mu^2) / 2 \partial \ln B / \partial t - (1 + \mu^2) (V / 2r) (1 - \cos^4 \psi)$$

$$\partial \ln B / \partial t = 0 \quad d \ln p / dt = -(1 + \mu^2) (V / 2r) (1 - \cos^4 \psi)$$

$$= -(1 + \mu^2) (V / 2r) [\sin^2 \psi + (1/4) \sin^2 2\psi]$$

Limit: $\tan \psi \ll 1$ or $r \sin \Theta \ll V / \Omega \sim 1 \text{ AU}$

$$d \ln p / dt \rightarrow -(V / 2r) (1 + \mu^2) 2 (\Omega r \sin \Theta / V)^2$$

$$= -(1 + \mu^2) (\Omega \sin \Theta)^2 (r / V) \rightarrow 0 \text{ as } r \rightarrow 0$$

No “adiabatic deceleration” near the Sun!

Conclusions for guiding center energy loss in a Parker magnetic field (applicable to “scatter-free” propagation)

Outside Earth’s orbit: $r \sin \Theta > V/\Omega \sim 1 \text{ AU}$

$\langle d \ln p / dt \rangle \rightarrow -2V/3r$ “adiabatic cooling”

Inside Earth’s orbit: $r \sin \Theta < V/\Omega \sim 1 \text{ AU}$

$$d \ln p / dt = (1 - \mu^2) / 2 \partial \ln B / \partial t - (1 + \mu^2) (V / 2r) (1 - \cos^4 \psi)$$

$$\rightarrow -(1 + \mu^2) (\Omega \sin \Theta)^2 (r / V) \rightarrow 0 \text{ as } r \rightarrow 0$$

**No “adiabatic deceleration” near the Sun!—for
any nearly “scatter-free” particle (PUIs or GCRs)**